

# After 2 years of PhD

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# Contents

- 1 Aspects of Quantum Fermionic T-duality
- 2 Supersymmetric Fluid Dynamics from Black Holes Superpartners
- 3 Conclusions

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# Review 1: Compactification

## Compactification of Closed Bosonic String

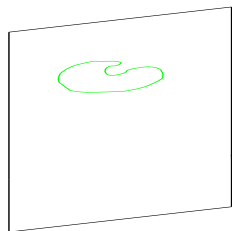
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Transformation  $W \leftrightarrow K$ ,  $R \leftrightarrow \frac{1}{R}$ :

- leaves  $M^2$  invariant
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⇒ Target Space Duality

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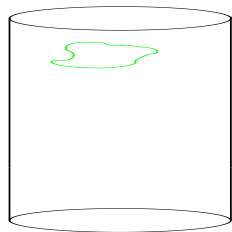
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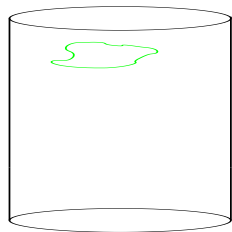
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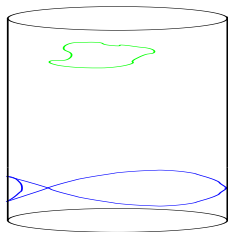
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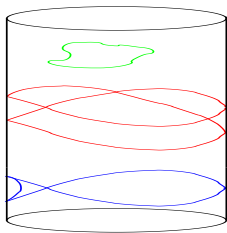
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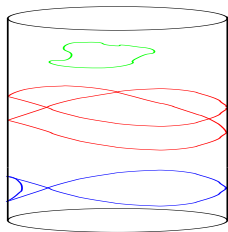
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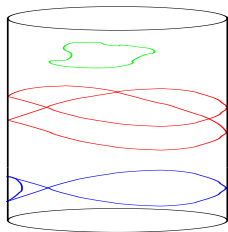
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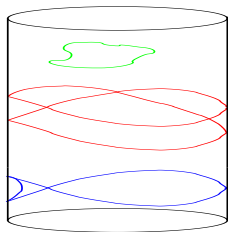
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$\Rightarrow$  Target Space Duality

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- action invariant under  $y \rightarrow y + \lambda$

$$S = \int d^2z \left( G_{ab}(x) \partial x^a \bar{\partial} x^b + G_{00}(x) \partial y \bar{\partial} y \right)$$

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- $G_{00} \rightarrow \frac{1}{G_{00}}$ : Target Space Transformation  $\rightarrow$  T-duality

## Review 3: Developments

- T-duality construction for **abelian isometries**  $x \rightarrow x + \lambda$

T.H. Buscher, *Phys. Lett.* **B 201** (1988)

- T-duality construction for **non-abelian, linear isometries**  $x \rightarrow x + \lambda x$

X.C. de la Ossa, F. Quevedo – hep-th/9210021

- Dual Superconformal Symmetry: **CFT counterpart of AdS Self-T-duality**

Self-T-duality of  $AdS_5 \times S^5$ : **bosonic  $\oplus$  fermionic T-duality**

N. Berkovits, J. Maldacena – arXiv:0807.3196

N. Beisert, R. Ricci, A.A. Tseytlin, M. Wolf – arXiv:0807.3228

- Limit of  $AdS_5 \times S^5$ : **purely fermionic coset  $\sigma$ -model**

N. Berkovits – hep-th/0703282

We Need an Analysis of Quantum Fermionic T-duality

# Outlook

## Playground:

- Coset spaces technique for  $\sigma$ -models
- $OSp(n|m)$  supergroup
- Examples and general case  $\frac{OSp(n|m)}{SO(n) \times Sp(m)}$

## Fermionic T-duality:

- Simple examples  $\longrightarrow$  Obstructions
- New methods
- Dualization of general  $OSp$  fermionic model

## Quantum Analysis:

- BFM method
- 1- and 2-loop self energy computation

# Fermionic Coset Spaces

## Generalities

- Coset group  $G/H \rightarrow$  algebra  $\mathfrak{g}$ , generators  $K^M = \left\{ T^a \in \frac{\mathfrak{g}}{\mathfrak{h}}, H_i \in \mathfrak{h} \right\}$
- Coset representative: choice of  $X_A$  in  $g = e^{K^A X_A}$
- Vielbein:  $V^a: g^{-1} dg = V^a T_a + \text{H-connection}$
- Sigma model action:  $S = \int_{\Sigma} d^2 z \sqrt{-\eta} \eta^{\mu\nu} V_{\mu}^a V_{\nu}^b \text{tr} [T_a T_b]$

## Fermionic Coset

- Bosonic part is always a subgroup of the isometry supergroup:

$$g = \exp \left( \sum_i x_i T_i + \sum_{\alpha} \theta_{\alpha} Q_{\alpha} \right) \rightarrow \exp \left( \sum_{\alpha} \theta_{\alpha} Q_{\alpha} \right)$$

- Few fermions:  $g \rightarrow \prod_{\alpha} \exp(\theta_{\alpha} Q_{\alpha})$

## Simple Examples

$$\frac{Osp(1|2)}{Sp(2)}:$$

- Action:

$$S \sim \int_{\Sigma} d^2z (1 + \theta_1\theta_2) \varepsilon_{\alpha\beta} \partial\theta^\alpha \bar{\partial}\theta^\beta$$

- Isometry:

$$\theta_i \rightarrow \theta_i + (1 + \theta_1\theta_2) \epsilon_i \quad \epsilon_i = \text{const}$$

How to perform T-duality?

- Action:

$$\frac{Osp(2|2)}{SO(2) \times Sp(2)}:$$

$$S \sim \int_{\Sigma} \left\{ d\theta_1 \wedge *d\theta_4 - d\theta_2 \wedge *d\theta_3 - 4\theta_3\theta_4 d\theta_1 \wedge *d\theta_2 \right\}$$

- Two translational isometries
- Equation of Motion:  $4\theta_3\theta_4 * A_2 = *d\theta_4 - d\tilde{\theta}_2$

Impossible to invert the EoM of A

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- general coset model:  $S = \int_{\Sigma} d^2z \sqrt{-\eta} \eta^{\mu\nu} V_{\mu}^a V_{\nu}^b \text{tr}[T_a T_b]$
- isometries generated by Killing vectors  $X^i \rightarrow X^i + \lambda^{\Lambda} K_{\Lambda}^i$
- transformation of the Vielbein:  $V^A \rightarrow V^A + d\lambda^{\Lambda} i_{K_{\Lambda}} V^A + \dots$
- gauging:  $V^a \rightarrow V^a + A^a$

gauge fixing  $\Rightarrow$  Solving

$$d\lambda^{\Lambda} = -V^a (i_{K_{\Lambda}} V^a)^{-1}$$

Two conditions:

- i existence of a minor:  $\exists M_I^S \equiv (i_{K_I} V^S)$  invertible
- ii existence of a solution of  $d[V^S (M^{-1})_S^g] = 0$

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# Non Linear Isometries

- Action for coset spaces models: depends on parametrization
- Non-linear isometry  $\rightarrow$  cumbersome procedure
- Fermionic coordinates: change of basis not always possible

## Wayout: Plücker Relations

- Introduction of new coordinates and constraints
- Solving the constraints  $\rightarrow$  back to the original action
- New action: the isometries act linearly

$$\mathcal{L} \sim (1 + \theta_1 \theta_2) \varepsilon_{\alpha\beta} \partial \theta^\alpha \bar{\partial} \theta^\beta \rightarrow \mathcal{L} \sim \varepsilon_{\alpha\beta} \partial \theta^\alpha \bar{\partial} \theta^\beta - \partial \phi \bar{\partial} \phi + \alpha (\phi^2 - \theta^2 - 1)$$

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- Lagrangian

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- Action invariant under the whole  $OSp(1|2)$

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- Collecting the gauge fields

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- Gauging:  $\nabla\theta^\alpha = \partial\theta^\alpha - A^\alpha \phi - A^\alpha{}_\beta \theta^\beta$ ,  $\nabla\phi = \partial\phi - A^\alpha \theta_\alpha$

- Collecting the gauge fields

$$\mathcal{L}_{gauging} \sim \mathcal{L}_0 + \left( h^\alpha + f^{\alpha\beta} A_\beta + g^{\alpha(\beta\gamma)} A_{\beta\gamma} \right) \bar{A}_\alpha + \left( l^{(\alpha\beta)} + m^{(\alpha\beta)\rho} A_\rho + n^{(\alpha\beta)(\rho\sigma)} A_{\rho\sigma} \right) \bar{A}_{\alpha\beta} + \bar{h}^\alpha A_\alpha + \bar{l}^{(\alpha\beta)} A_{\alpha\beta}$$

# New Method: $\frac{OSp(1|2)}{Sp(2)}$

- Lagrangian

$$\mathcal{L}_0 \sim \varepsilon_{\alpha\beta} \partial\theta^\alpha \bar{\partial}\theta^\beta - \partial\phi \bar{\partial}\phi + \alpha (\phi^2 - \theta^2 - 1)$$

- Action invariant under the whole  $OSp(1|2)$

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- Gauge fixing  $\oplus$  EoM for gauge fields  $\rightarrow$  T-dual model

# General Model: $\frac{OSp(n|m)}{SO(n) \times Sp(m)}$

- Supergroup representative:  $L = \begin{pmatrix} \Lambda_j^i & \Theta_\alpha^i \\ \Theta_{\alpha j} & \Phi_{\beta}^{\alpha} \end{pmatrix}$

- $\Lambda_{(ij)}$  bosonic  $SO(n)$  fields;
- $\Phi_{[\alpha\beta]}$  antisymmetric  $Sp(m)$  fields;
- $\Theta_{i\alpha}$  fermionic fields.

- Vielbein: expand  $L^{-1}\partial L$  over  $\frac{osp(n|m)}{so(n) \times sp(m)}$  generators

$$\mathcal{L} = \mathcal{L}_0 + \mathcal{L}_P$$

- $\mathcal{L}_V = V^\alpha_i \delta^{ij} \varepsilon_{\alpha\beta} V^\beta_j + V^i_\alpha \delta_{ij} \varepsilon^{\alpha\beta} V^j_\beta$

- $\mathcal{L}_P = \alpha^{(ij)} \left( \Lambda_{ik} \delta^{kl} \Lambda_{lj} - \Theta_{i\alpha} \varepsilon^{\alpha\beta} \Theta_{j\beta} - \delta_{ij} \right) +$   
 $\beta^{[\alpha\beta]} \left( \Phi_{\alpha\gamma} \varepsilon^{\gamma\delta} \Phi_{\delta\beta} - \Theta_{i\alpha} \delta^{ij} \Theta_{j\beta} - \varepsilon_{\alpha\beta} \right) + \gamma^{i\alpha} \left( \Lambda_{ik} \delta^{kl} \Theta_{k\alpha} + \Phi_{\alpha\beta} \varepsilon^{\beta\gamma} \Theta_{i\gamma} \right)$

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Construction of T-dual Model!

# $\frac{OSp(n|m)}{SO(n) \times Sp(m)}$ : Quantum Analysis

- Computation of 1- and 2-loop correction to self energy

## Background Field Method:

- Group Representative:  $g = g_0 e^{\lambda X}$
- background external field  $B_\mu = g_0^{-1} \partial_\mu g$ , quantum internal field  $X$
- Vielbein:  $V_\mu = B_\mu + \frac{\lambda^2}{2} [[B_\mu, X], X] + \lambda \partial_\mu X + \frac{\lambda^3}{3!} [[\partial_\mu X, X], X] + \dots$
- Action:  $S = \frac{1}{2\pi\lambda^2} \int \text{Str}(V \cdot V) \longrightarrow$  Feynman Diagrams

Self Energy, both 1 and 2-loops  $\propto (2 + m - n)$

- 1-loop computation for T-dual model  $\longrightarrow$  same result!

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# Results and Open Issues

## Classical Analysis:

- Study of obstructions to Fermionic T-duality
- New method based on Plücker relations
- Construction of T-dual of  $\frac{OSp(n|m)}{SO(n) \times Sp(m)}$

## Quantum Analysis:

- BFM
- 1- and 2-loop Self Energy
- 1- and 2-loop Self Energy for T-dual model

## Open Problems:

- T-duality for other models
- Analysis of WZW models
- Deep understanding of quantum T-duality

# Contents

- 1 Aspects of Quantum Fermionic T-duality
- 2 Supersymmetric Fluid Dynamics from Black Holes Superpartners**
- 3 Conclusions

# Fluid-Gravity Correspondence 1

- *AdS/CFT* correspondence

Type IIB string theory on asymptotically  
 $AdS_5 \times S^5$

$\mathcal{N} = 4$  SYM on  $S^3 \times R^1$  or  $R^3 \times R^1$

- Effective Description: length scales  $\gg$  mean free path length

Einstein Equations with negative  
cosmological constant

Relativistic Fluid Dynamics

- Static Solution

Black Hole or Black Branes in *AdS*...

static perfect fluid configuration

- Perturbation

Evolving Black Brane

dissipative fluid flow

# Fluid-Gravity Correspondence 2

## Interesting Recent Results from Fluid-Gravity Correspondence:

Battacharyya, Hubeney, Minwalla, Rangamani – arXiv:0712.2456

- Perturbative procedure for Fluid-Gravity Correspondence
- Gravitational derivation of Navier-Stokes Equations

fluid dynamic = long wavelength effective theory for strongly coupled QFT

- Explicit construction of Fluid Energy Momentum Tensor

new method to compute Transport Coefficients

# Basic Setup

- Sector of IIB SUGRA: pure gravitational dynamics in asymptotically  $AdS$
- $AdS_5$  Black Brane in Poincaré patch:

$$2ds_{BH} = dv dr - r^2 f(r) dv^2 + r^2 dx^2$$

$$f(r) = 1 - \frac{1}{r^4}$$

- Subgroup of  $SO(2, 4) \rightarrow$  Boosted Black Brane

$$d^2s = -2u_\mu dx^\mu dr + r^2 (1 - f(br)) u_\mu u_\nu dx^\mu dx^\nu + r^2 \eta_{\mu\nu} dx^\mu dx^\nu$$

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# Perturbative Expansion

- Parameters functions of boundary coords  $\rightarrow$  not a solution of Einstein Eq.
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$$G = G^0(\beta_i, b) + \varepsilon G^1(\beta_i, b) + O(\varepsilon^2)$$

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$$\partial_\nu \mathbf{b}^{(0)} = \frac{1}{3} \partial_i \beta_i^{(0)}, \quad \partial_i \mathbf{b}^{(0)} = \partial_\nu \beta_i^{(0)}$$

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- Pure Gravity  $g_{\mu\nu}$   $\longrightarrow$  Pure Supergravity  $g_{\mu\nu}, \psi_\mu$ :
- Supersymmetry transformation:

$$\delta_\epsilon g_{\mu\nu} = \text{Re}(\bar{\epsilon} \Gamma_{(\mu} \psi_{\nu)}) , \quad \delta_\epsilon \psi_\mu = \left( D_\mu(\omega) + \frac{1}{2} \Gamma_\mu \right) \epsilon$$

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Burrington, Liu, Sabra: hep-th/0412155

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# Killing Vectors

- Bosonic Background:  $AdS_5$  with boundary interpolating coordinates

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- Computation of the 8  $AdS_5$  Killing Spinors  $\epsilon$ :

$$\delta_\epsilon \psi_\mu = \left( D_\mu (\omega) + \frac{1}{2} \Gamma_\mu \right) \epsilon = 0$$

- Notice that  $\bar{\epsilon} \Gamma^\mu \epsilon \partial_\mu$  is the  $AdS_5$  Killing vector
- Turning on the BH:  $\delta_\epsilon \psi_\mu \neq 0$
- Supersymmetry  $\rightarrow$  BH Superpartner:

$$\delta_\epsilon^2 g_{\mu\nu} = \text{Re} \left( \bar{\epsilon} \Gamma_{(\mu} \delta \psi_{\nu)} \right)$$

- New metric term depending on fermionic bilinears  $\{\lambda, l_i\}$

# Results

- We reproduce the computation by Minwalla *et al.*
- 1<sup>st</sup>-order metric with **local** parameters

$$g = g^{BH} + \varepsilon \left[ \mathcal{L}_K g_{BH}(\beta_i, b, \alpha_i, a) + \delta_\varepsilon^2 g_{BH}(l_i, \lambda) + g^1 \right]$$

- Imposing Einstein Equations in large  $r$  limit  $\rightarrow$  4 New Constraints eq's:

$$(\partial_i \beta_i - 3\partial_0 b) + (\partial_i l_i + \partial_0 \lambda) = 0, \quad (\partial_0 \beta_i - \partial_i b) - (\partial_i \lambda + 3\partial_0 l_i) = 0$$

- Correction to NS equations!

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# Road to Supersymmetric Fluid Dynamics 1

Interpretation of modified NS?  $\longrightarrow$  we need a **Supersymmetric Fluid Dynamics**

P.A. Grassi, AM, L. Sommovigo – arXiv:1107.2780

- Idea: supersymmetrize an appropriate bosonic Lagrangian
- Bosonic model using **Clebsh Parametrization**:

$$\mathcal{L} = \sqrt{-g} \left( j^\mu (\partial_\mu a + \alpha \partial_\mu \beta) + f(j^2) \right)$$

- EoM for  $a \longrightarrow \partial_\mu j^\mu = 0$
- Other EoM's  $\oplus$  manipulations  $\longrightarrow$  NS equations (no dissipation terms)

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## Supersymmetrization

- $j_\mu \rightarrow$  linear superfield  $J$  satisfying  $\bar{D}DJ = 0$
- $\partial_\mu a + \alpha \partial_\mu \beta \rightarrow$  real superfield  $A$  in Wess-Zumino gauge
- introduction of  $\mathcal{J}^\mu \propto (\bar{D}\gamma_5\gamma_\mu D) J$

$$S = \int d^4x \int d^4\theta \left( -JA + F(\mathcal{J}_\mu \mathcal{J}^\mu) \mathcal{J}^2 \right)$$

## Results

- Construction of Supersymmetric Lagrangian
- Computation of the complete component expansion  $\rightarrow$  different sector analysis
- Supersymmetric version of Clebsh parametrization ( not seen here )



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# Results and Open Issues

## Fluid-Gravity Correspondence:

- Application of perturbative procedure *a la* Minwalla to Black Hole Superpartners
- Derivation of NS equations modified by fermionic dof's

## Models for Supersymmetric Fluid Dynamics:

- Construction of suitable supersymmetric lagrangian
- Analysis of different sectors

## In preparation:

- Computation of boundary Energy Momentum Tensor
- Derivation of modified NS from supersymmetric Lagrangian
- Supergravity Extension

# Contents

- 1 Aspects of Quantum Fermionic T-duality
- 2 Supersymmetric Fluid Dynamics from Black Holes Superpartners
- 3 Conclusions

# Conclusions

## List of Publications

- Aspects of Quantum Fermionic T-duality

P.A. Grassi, AM – arXiv:1101.5969

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Thank You!

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